

# Sterile neutrinos from D-brane models

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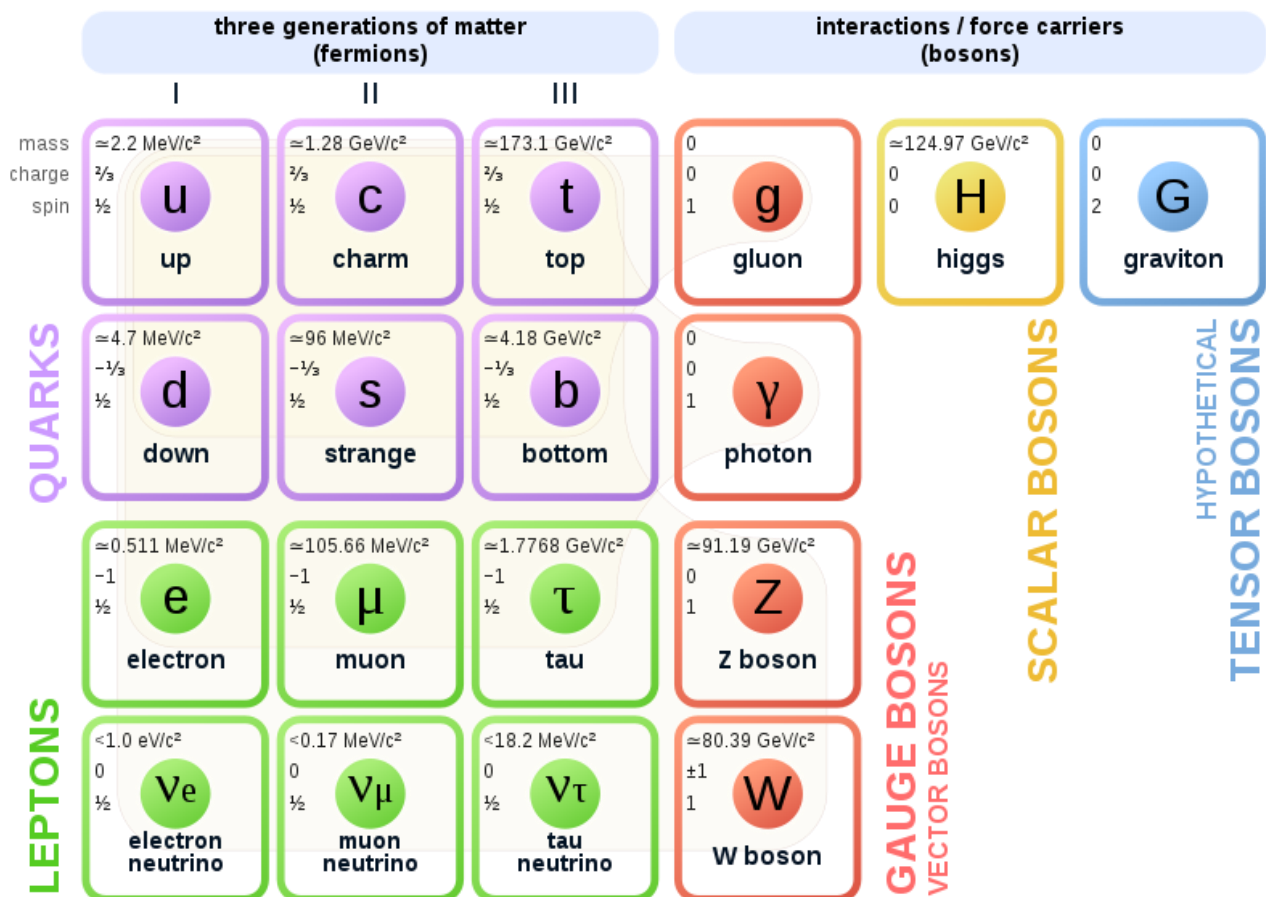
21<sup>st</sup> String Phenomenology – 2022  
Southampton

## OUTLINE

- Standard model “weaknesses”
- Neutrino mass in the SM & beyond/Strings
- ~ Intersecting D-brane models :
  - ✓ Gauged baryon number (stable proton)
    - ✓ Left + Right handed neutrinos
    - + Sterile neutrinos

# → STANDARD MODEL

## Standard Model of Elementary Particles and Gravity



●+ accommodates 3 generations of neutrinos

# BUT : “weaknesses”

- Baryon **non**-conservation  
(B-L conserved)

⇒ Proton decay

⇒

$$O_{abcd}^{(1)} = (\bar{d}_{\alpha a R}^C u_{\beta b R}) (\bar{q}_{i \gamma c L}^C l_{j d L}) \epsilon_{\alpha \beta \gamma} \epsilon_{i j},$$

$$O_{abcd}^{(2)} = (\bar{q}_{i \alpha a L}^C q_{j \beta b L}) (\bar{u}_{\gamma c R}^C l_{d R}) \epsilon_{\alpha \beta \gamma} \epsilon_{i j},$$

$$O_{abcd}^{(3)} = (\bar{q}_{i \alpha a L}^C q_{j \beta b L}) (\bar{q}_{k \gamma c L}^C l_{l d L}) \epsilon_{\alpha \beta \gamma} \epsilon_{i j} \epsilon_{kl},$$

$$O_{abcd}^{(4)} = (\bar{q}_{i \alpha a L}^C q_{j \beta b L}) (\bar{q}_{k \gamma c L}^C l_{l d L}) \epsilon_{\alpha \beta \gamma} \\ \times (\vec{\tau} \epsilon)_{i j} \cdot (\vec{\tau} \epsilon)_{kl},$$

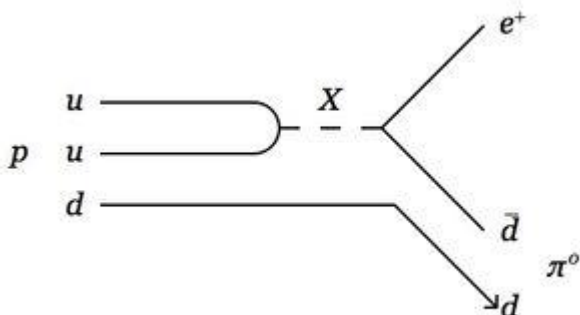
$$O_{abcd}^{(5)} = (\bar{d}_{\alpha a R}^C u_{\beta b R}) (\bar{u}_{\gamma c R}^C l_{d R}) \epsilon_{\alpha \beta \gamma},$$

$$O_{abcd}^{(6)} = (\bar{u}_{\alpha a R}^C u_{\beta b R}) (\bar{d}_{\gamma c R}^C l_{d R}) \epsilon_{\alpha \beta \gamma}.$$

Weinberg (1979)

Wilczek and Zee (1979)

- $p \rightarrow e^+ \pi^0$   $\tau_p > 8.2 \times 10^{33}$  years  $\Lambda \gtrsim 10^{16}$  GeV



Hierarchy problem 100 GeV –  $10^{16}$  GeV

- **No Gravity**

- # No Mass term for neutrinos

■ Motivation for introducing neutrino mass term → Beyond the SM

- A mass can only be introduced beyond the SM e.g. by adding a right handed neutrino (s) / (see saw mechanism)

- **Neutrinos have a mass** - Discovery of **neutrino oscillations**

- Explain (inconclusive) excess of low energy electronic recoil events, over known backgrounds, observed at XENON1T experiment at 2-3 KeV ?

- **Dark Matter** (DM) contributes five times more to the energy of the Universe than ordinary matter (Weakly interacting) dark matter candidates => sterile neutrinos of KeV masses + with small mixing with active neutrinos.

Light sterile neutrinos : A white paper,  
<https://arxiv.org/pdf/1204.5379.pdf>;

- T2K experiment → **STERILE NEUTRINOS  $< 1$  eV**

# ? Solution – Building a model BEYOND the SM without Proton decay

We need to find the particle content of the :  
**Standard Model** .The heaviest elementary particles  
on the right side ...

Three Generations of Matter (Fermions)

	I	II	III	
mass→	2.4 MeV	1.27 GeV	171.2 GeV	0
charge→	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0
spin→	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
name→	u up	c charm	t top	$\gamma$ photon
	4.8 MeV	104 MeV	4.2 GeV	0
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
	d down	s strange	b bottom	g gluon
	$< 2.2$ eV	$< 0.17$ MeV	$< 15.5$ MeV	91.2 GeV
	0	0	0	0
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
	$\nu_e$ electron neutrino	$\nu_\mu$ muon neutrino	$\nu_\tau$ tau neutrino	Z weak force
	0.511 MeV	105.7 MeV	1.777 GeV	80.4 GeV
	-1	-1	-1	$\pm 1$
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
	e electron	$\mu$ muon	$\tau$ tau	W <sup>±</sup> weak force

Quarks

Leptons

Bosons (Forces)

Build

**STRING THEORY**

**Standard model-like model**

1 Macroscopic level - matter, 2 Molecular level,  
3 Atomic level, 4 Electrons, 5 Quarks, 6 String Theory



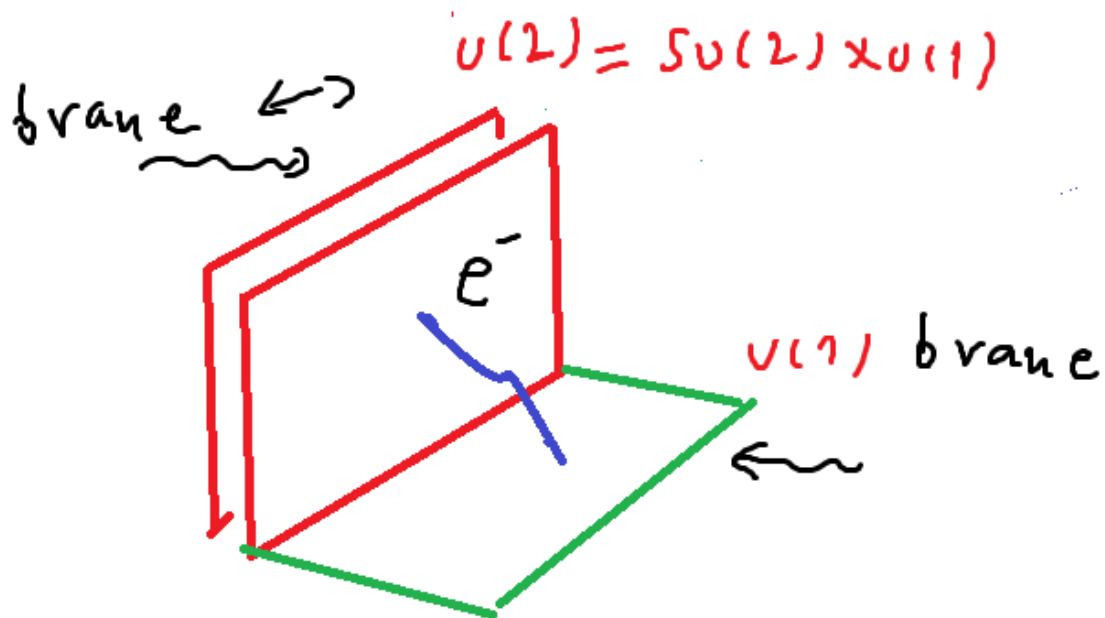


Particles  $\rightarrow$  localized among intersecting branes

What is an intersecting brane ?

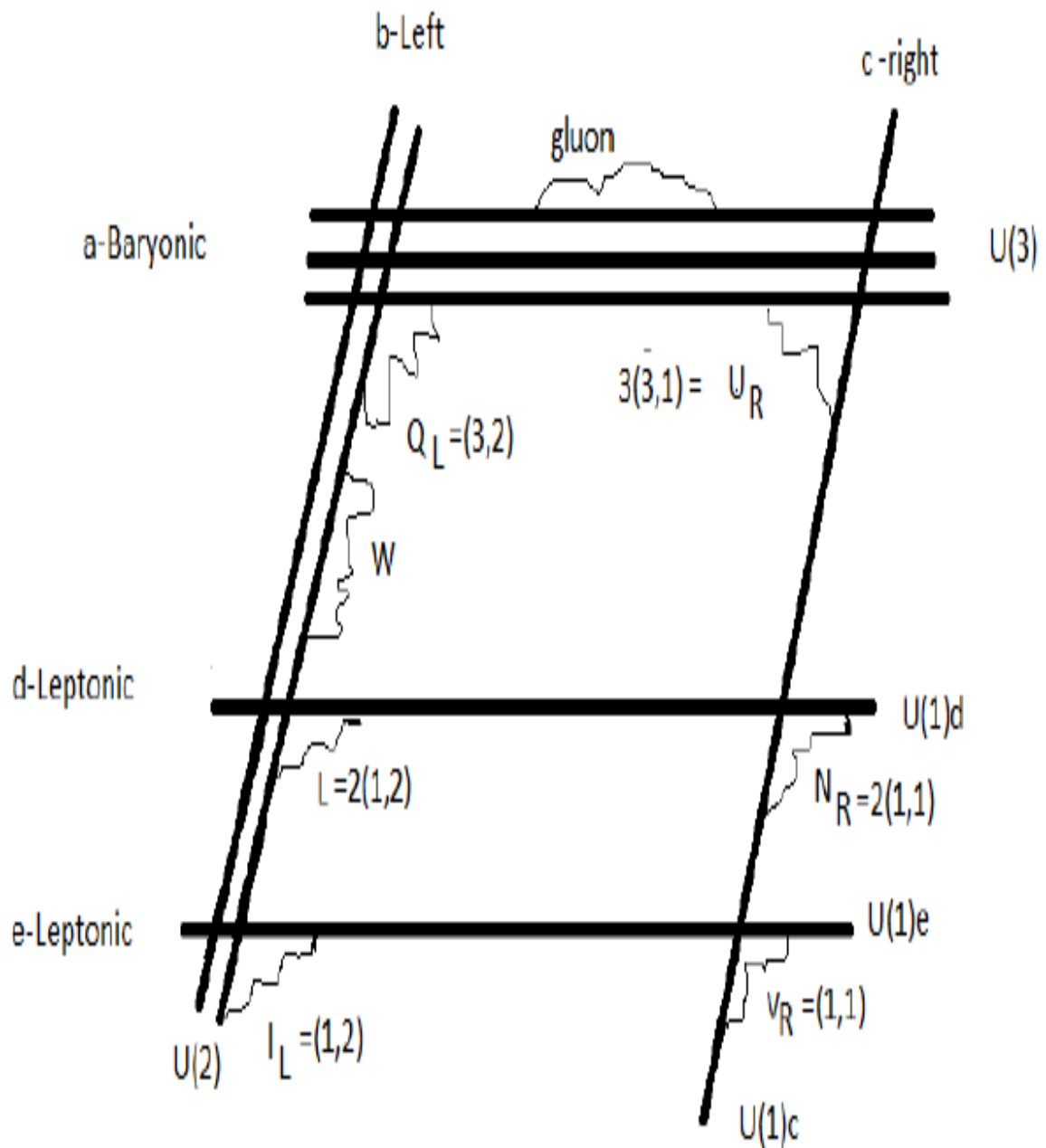
A higher dimensional hypersurface

Simplest representation



$e^- \rightarrow (2,1)$  representation  $\rightarrow$  and charges  $(1, -1) = (Q_a, Q_b)$

# 5-stack String Standard Model



SU(3)<sub>a</sub> SU(2)<sub>b</sub> U(1)<sub>a</sub> U(1)<sub>b</sub> U(1)<sub>c</sub> U(1)<sub>d</sub> U(1)<sub>e</sub>

C.K (2002)

# MATTER SPECTRUM

Matter Fields		Intersection	$Q_a$	$Q_b$	$Q_c$	$Q_d$	$Q_e$	Y
$Q_L$	$(3, 2)$	$I_{ab} = 1$	1	-1	0	0	0	1/6
$q_L$	$2(3, 2)$	$I_{ab^*} = 2$	1	1	0	0	0	1/6
$U_R$	$3(\bar{3}, 1)$	$I_{ac} = -3$	-1	0	1	0	0	-2/3
$D_R$	$3(\bar{3}, 1)$	$I_{ac^*} = -3$	-1	0	-1	0	0	1/3
$L$	$2(1, 2)$	$I_{bd} = -2$	0	-1	0	1	0	-1/2
$l_L$	$(1, 2)$	$I_{be} = -1$	0	-1	0	0	1	-1/2
$N_R$	$2(1, 1)$	$I_{cd} = 2$	0	0	1	-1	0	0
$E_R$	$2(1, 1)$	$I_{cd^*} = -2$	0	0	-1	-1	0	1
$\nu_R$	$(1, 1)$	$I_{ce} = 1$	0	0	1	0	-1	0
$e_R$	$(1, 1)$	$I_{ce^*} = -1$	0	0	-1	0	-1	1

Table 1: Low energy fermionic spectrum of the five stack string scale  $SU(3)_C \otimes SU(2)_L \otimes U(1)_a \otimes U(1)_b \otimes U(1)_c \otimes U(1)_d \otimes U(1)_e$ , type I D6-brane model together with its  $U(1)$  charges. Note that at low energies only the SM gauge group  $SU(3) \otimes SU(2)_L \otimes U(1)_Y$  survives.

$$Q_a = 3B, \quad L = Q_d + Q_e$$

“Predicts...”

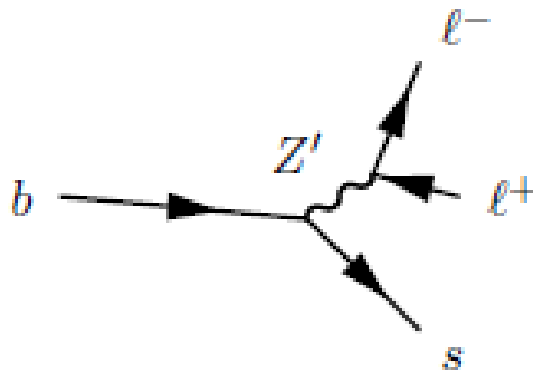
- the existence of 1 or 2 sneutrinos
- > break the extra  $Z'$  with  $\nu_R$

=> Used...  $\nu_R$  in MSSM to break  $U(1)_{B-L}$   
Barger, Perez, Spineer (2009)

- explains LHCb b-anomalies

$$R_K = \frac{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)}{dq^2} dq^2}{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\Gamma(B^+ \rightarrow K^+ e^+ e^-)}{dq^2} dq^2}$$

deficit of muonic decays relative to electronic decays



A. Celis, W. Feng, D. Lust (2016)



Stringy Z' boson -> nonnegligible couplings to the first two quark generations

Z' Mass predicted  $\rightarrow \sim [3.5, 5.5]$  TeV,

should be possible to discover such a state directly during the next LHC runs via

Drell-Yan production in :

di-electron or

di-muon decay channels

$$\text{Br}(Z' \rightarrow \mu^+ \mu^-) / \text{Br}(Z' \rightarrow e^+ e^-) \sim [0.5-0.9]$$

# NEUTRINO MASSES



can originate via chiral symmetry breaking

C.K (2002)

Ibanez, Marchesano, Rabadan (2001)

$$\alpha'(LN_R)(Q_L U_R)^*, \quad \alpha'(l\nu_R)(q_L U_R)$$

From u-quark chiral condensate

$$\frac{\langle u_R u_L \rangle}{M_s^2} = \frac{(240 \text{ MeV})^3}{M_s^2}$$

$$M_\nu \sim (0.1-10 \text{ eV})$$

# STERILE NEUTRINOS

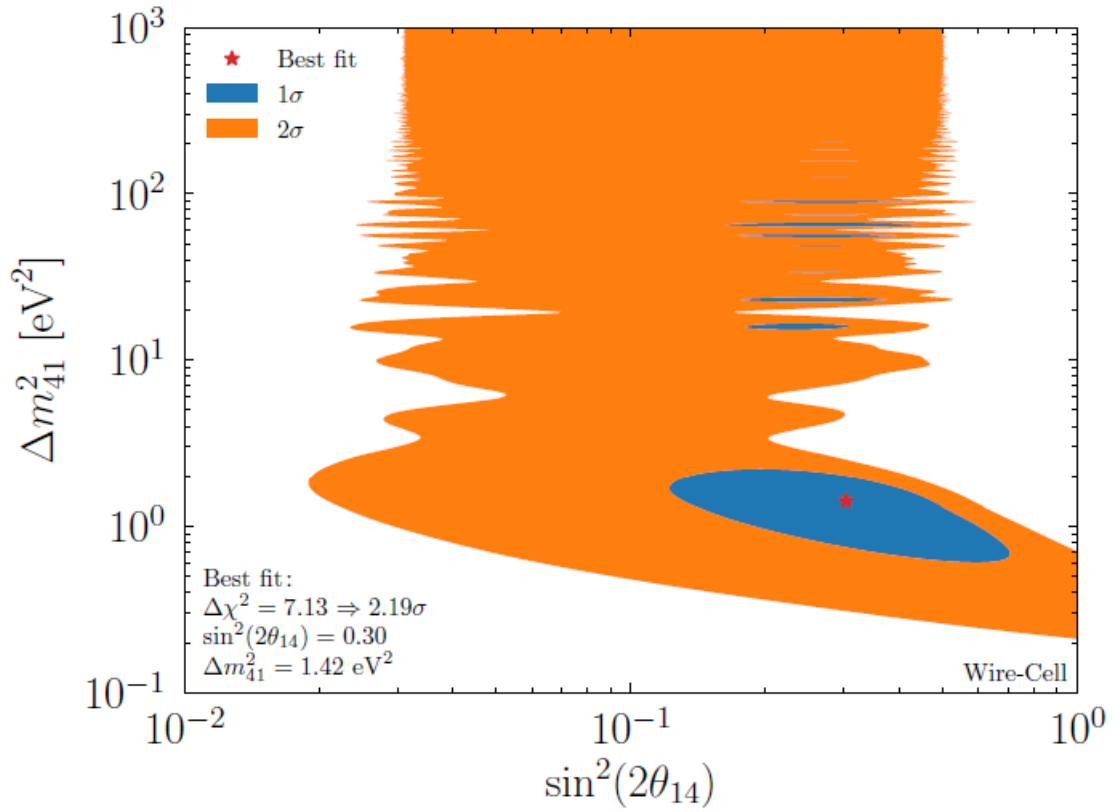


FIG. 2. The preferred regions in  $\Delta m_{41}^2$  -  $\sin^2(2\theta_{14})$  parameter space using data from MicroBooNE's Wire-Cell analysis [41]. The blue (orange) region is the preferred region at  $1\sigma$  ( $2\sigma$ ) assuming Wilks' theorem. The red star is at the best fit point:  $\Delta m_{41}^2 = 1.42 \text{ eV}^2$  and  $\sin^2(2\theta_{14}) = 0.30$  which has a test statistic of  $\Delta\chi^2 = 7.13$  to no oscillations which implies  $2.19\sigma$  under Wilks' theorem.

2111.05793[hep-ph]

- Sterile neutrinos in Gauge Theory

Inverse See Saw mechanism

$$\lambda_1 \nu_R \nu_L H + \lambda_2 \nu_R H N + \lambda_3 \frac{1}{M_{GUT}} \bar{K}^2 N N$$

$$m_D = \lambda_1 \langle H \rangle, \quad V_R = \lambda_2 \langle H \rangle, \quad \mu = \frac{\lambda_3}{M_{GUT}} \langle \tilde{K} \rangle^2$$

$$\begin{pmatrix} 0 & m_D^T & 0 \\ m_D & 0 & V_R \\ 0 & V_R & \mu \end{pmatrix}$$

Valle; Leontaris and Shafi

# ●● Sterile neutrinos in INTERSECTING D-BRANE models

$$\mathcal{L} = m_D \nu_L \nu_R + m_N \nu_R N_1 + m_\Sigma \nu_L N_1 + \dots$$

- Sterile neutrinos in eigenstate basis  
( $\nu_L, \nu_R, N_1$ )

- $\rightarrow$  mass matrix

$$\begin{pmatrix} 0 & m_D & 0 \\ m_D & 0 & m_N \\ 0 & m_N & 0 \end{pmatrix} \quad \begin{pmatrix} 0 & m_D & m_\Sigma \\ m_D & 0 & m_N \\ m_\Sigma & m_N & 0 \end{pmatrix}$$

**BARYON # CONSERVED**

I. Antoniadis and C.K

sterile neutrinos in other formulations  
(baryon # non-conserved  $\rightarrow$   
[Calabi-Yau, Fermionic, F-theory

Mohapatra & Valle (1986)  
,Faraggi & Guzzi; Leontaris]